

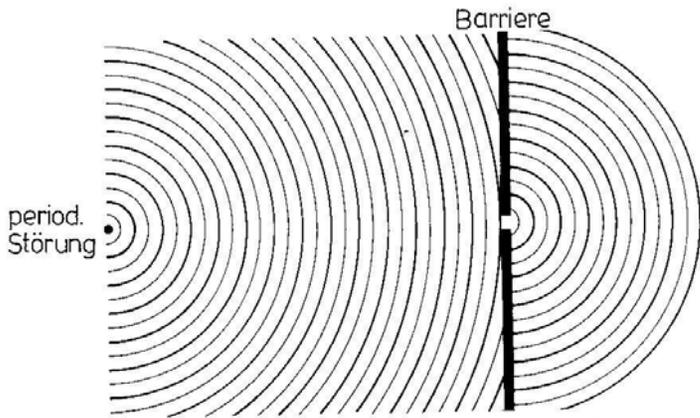
Elektrochemische Spannungsreihe

Tab. 2.8-2. Standard-Elektrodenpotentiale nach R. C. Weast (Hrsg.) Handbook of Chemistry and Physics, 60th Ed., 3. Nachdruck 1981, The Chemical Rubber Co, Boca Raton, Florida.

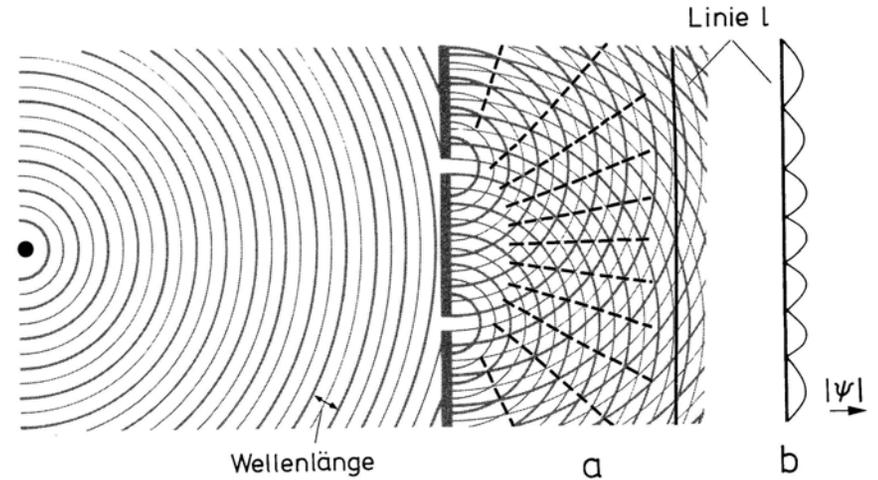
Halbzelle	Elektrodenreaktion	$\frac{E_h^0}{V}$
Metallionenelektroden		
Li ⁺ Li	Li ⁺ + e ⁻ ↔ Li	- 3.045
Rb ⁺ Rb	Rb ⁺ + e ⁻ ↔ Rb	- 2.925
K ⁺ K	K ⁺ + e ⁻ ↔ K	- 2.924
Cs ⁺ Cs	Cs ⁺ + e ⁻ ↔ Cs	- 2.923
Ca ²⁺ Ca	Ca ²⁺ + 2e ⁻ ↔ Ca	- 2.76
Na ⁺ Na	Na ⁺ + e ⁻ ↔ Na	- 2.7109
Mg ²⁺ Mg	Mg ²⁺ + 2e ⁻ ↔ Mg	- 2.375
Al ³⁺ Al	Al ³⁺ + 3e ⁻ ↔ Al	- 1.66
Zn ²⁺ Zn	Zn ²⁺ + 2e ⁻ ↔ Zn	- 0.7628
Fe ²⁺ Fe	Fe ²⁺ + 2e ⁻ ↔ Fe	- 0.409
Cd ²⁺ Cd	Cd ²⁺ + 2e ⁻ ↔ Cd	- 0.4026
Ni ²⁺ Ni	Ni ²⁺ + 2e ⁻ ↔ Ni	- 0.23
Pb ²⁺ Pb	Pb ²⁺ + 2e ⁻ ↔ Pb	- 0.1263
Cu ²⁺ Cu	Cu ²⁺ + 2e ⁻ ↔ Cu	+ 0.3402
Ag ⁺ Ag	Ag ⁺ + e ⁻ ↔ Ag	+ 0.7996
Au ⁺ Au	Au ⁺ + e ⁻ ↔ Au	+ 1.68
Gaselektroden		
H ⁺ H ₂ , Pt	2H ⁺ + 2e ⁻ ↔ H ₂	0.00
OH ⁻ O ₂ , Pt	O ₂ + 2H ₂ O + 4e ⁻ ↔ 4OH ⁻	+ 0.401
I ⁻ I ₂ , Pt	I ₂ + 2e ⁻ ↔ 2I ⁻	+ 0.535
Cl ⁻ Cl ₂ , Pt	Cl ₂ + 2e ⁻ ↔ 2Cl ⁻	+ 1.3583
F ⁻ F ₂ , Pt	F ₂ + 2e ⁻ ↔ 2F ⁻	+ 2.87
Elektroden 2. Art		
SO ₄ ²⁻ PbSO ₄ Pb	PbSO ₄ + 2e ⁻ ↔ Pb + SO ₄ ²⁻	- 0.356
I ⁻ AgI Ag	AgI + e ⁻ ↔ Ag + I ⁻	- 0.1519
Cl ⁻ AgCl Ag	AgCl + e ⁻ ↔ Ag + Cl ⁻	+ 0.2223
Cl ⁻ Hg ₂ Cl ₂ Hg	Hg ₂ Cl ₂ + 2e ⁻ ↔ 2Hg + 2Cl ⁻	+ 0.2682
Redoxelektroden		
Cr ³⁺ , Cr ²⁺ Pt	Cr ³⁺ + e ⁻ ↔ Cr ²⁺	- 0.41
Fe ³⁺ , Fe ²⁺ Pt	Fe ³⁺ + e ⁻ ↔ Fe ²⁺	+ 0.770
Chinhydron Pt	$ \text{O}=\text{C}_6\text{H}_4=\text{O} + 2\text{H}^+ + 2\text{e}^- \rightleftharpoons \text{HO}-\text{C}_6\text{H}_4-\text{OH} $	+ 0.6992
Ce ⁴⁺ , Ce ³⁺ Pt	Ce ⁴⁺ + e ⁻ ↔ Ce ³⁺	+ 1.4430

Wellennatur des Lichts – Analogie zu Wasserwellen

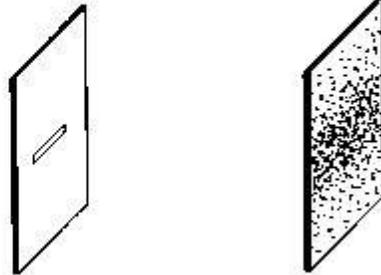
Beugung am Spalt



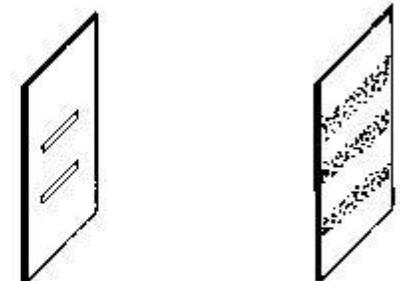
Interferenz am Doppelspalt



*
Lichtquelle
(monochromatisch,
punktförmig)

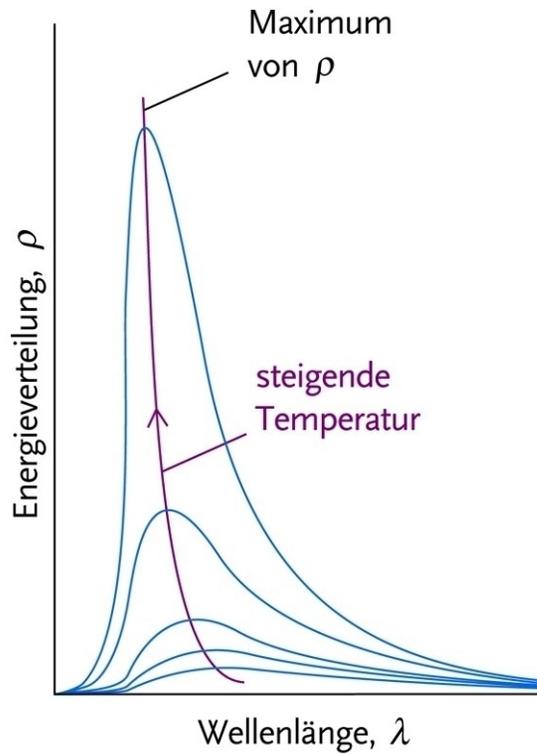


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Lichtquelle
(monochromatisch,
punktförmig)

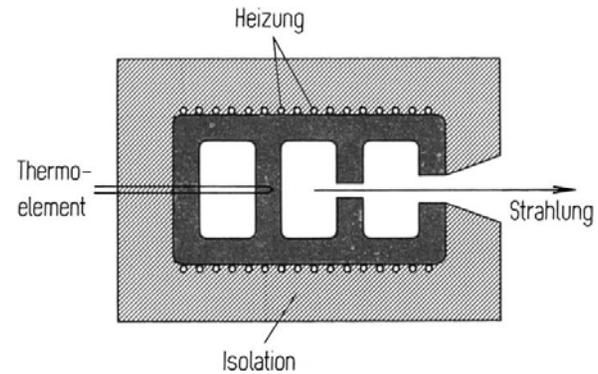
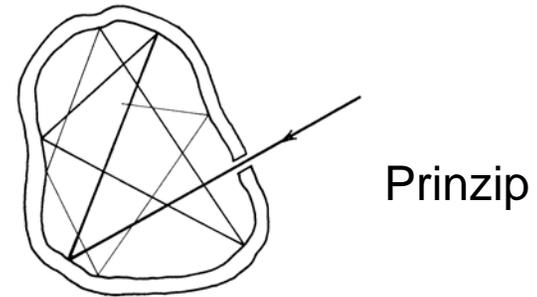


Energieverteilung eines schwarzen Strahlers

Typisches Spektrum

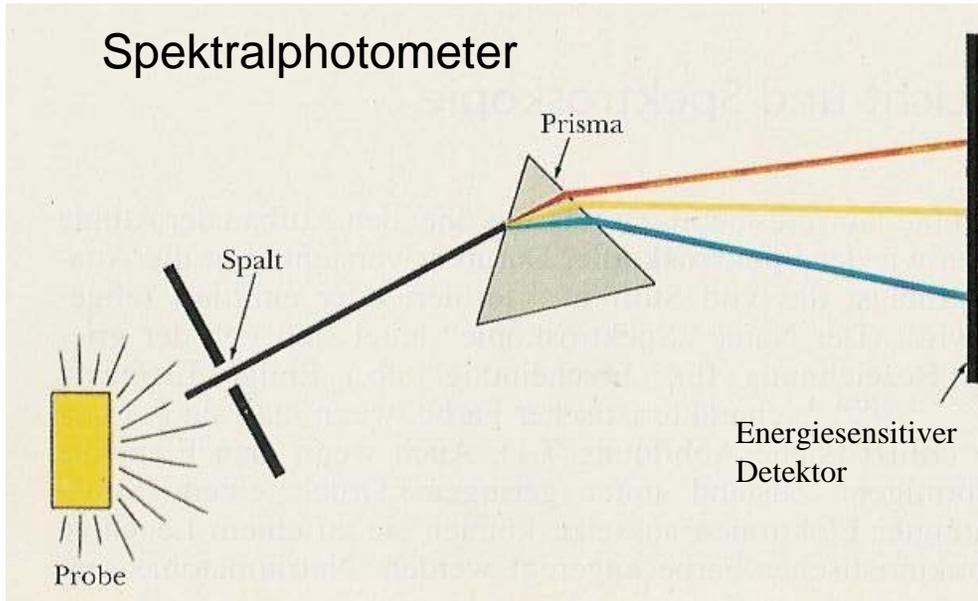


Experimentelle Realisierung eines Hohlraumstrahlers = Schwarzer Strahler



Experimenteller Aufbau

Messung der Schwarzkörper-Strahlung

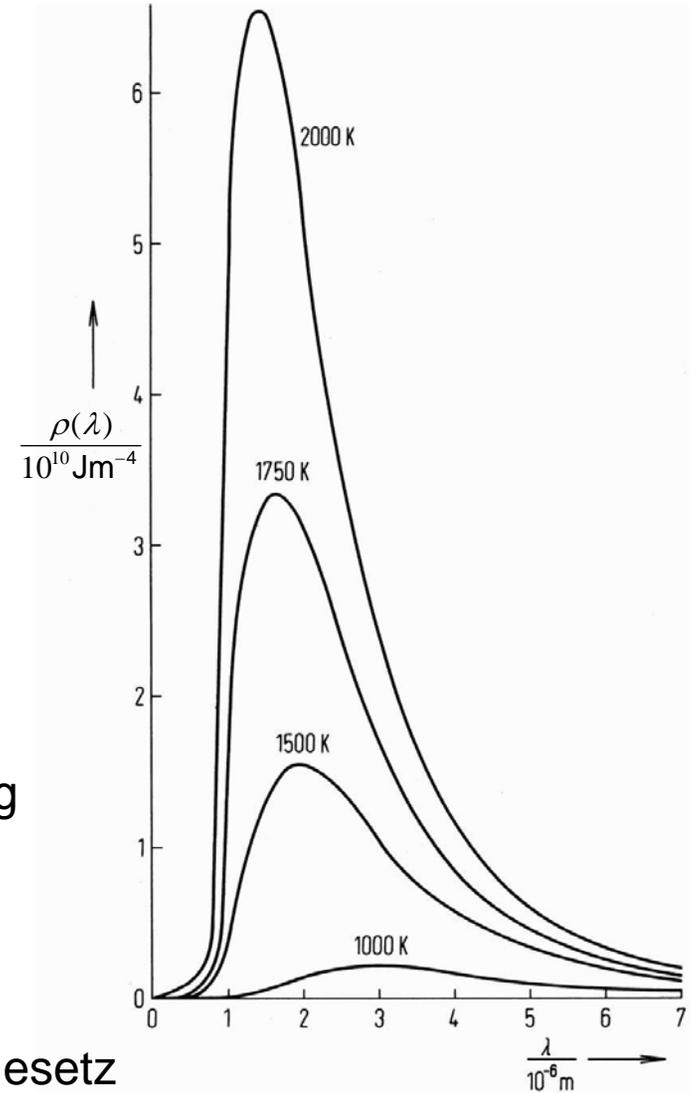


Zusammenhang zur Messgröße:

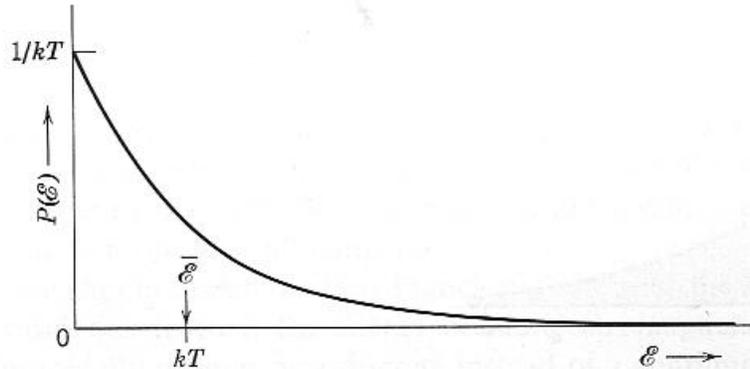
$$M = \frac{dQ}{dA dt} = \frac{1}{4} \rho c \quad \text{Spezifische Ausstrahlung}$$

$$M = \sigma \cdot T^4 \quad \text{Stefansches Gesetz}$$

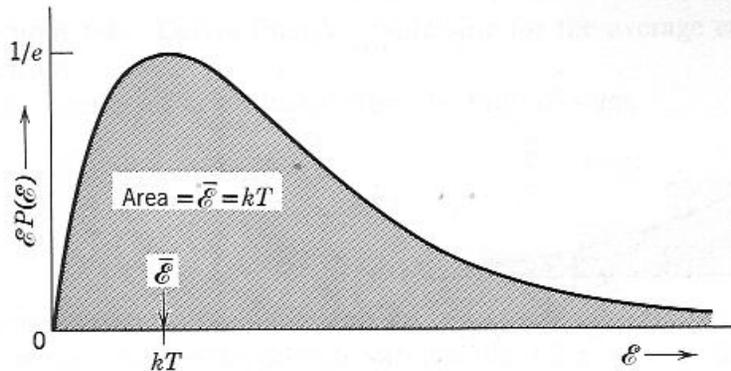
$$\lambda_{\max} \cdot T = \text{const} \quad \text{Wiensches Verschiebungs-Gesetz}$$



Boltzmannverteilung – Kontinuierlich



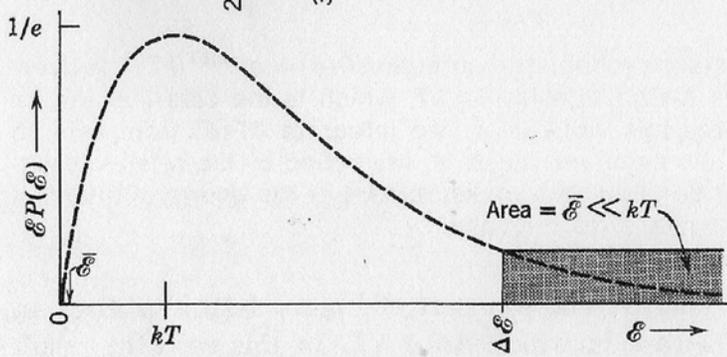
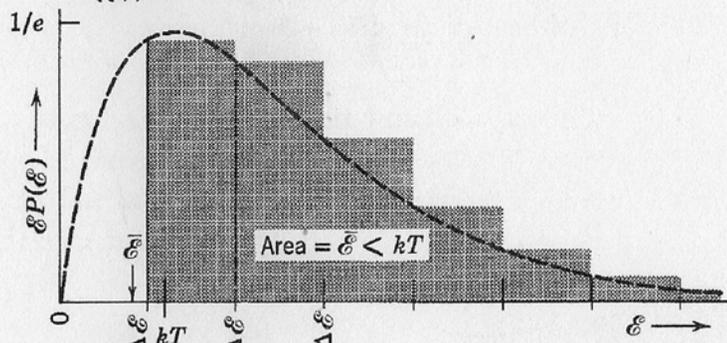
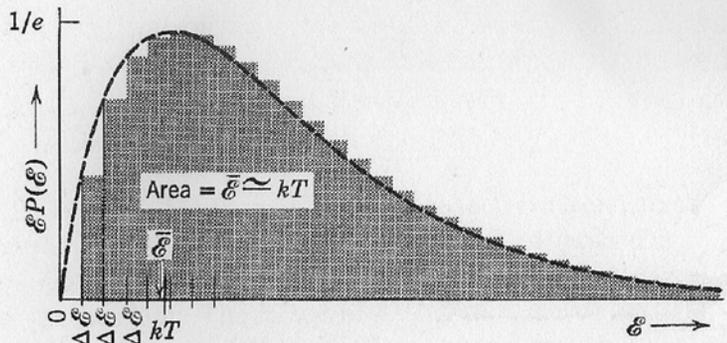
$$P(E) = \frac{1}{k_B T} e^{-E/k_B T}$$



$$\bar{E} = \frac{\int_0^{\infty} E \cdot P(E) dE}{\int_0^{\infty} P(E) dE} = k_B T$$

Figure 1-9 *Top:* A plot of the Boltzmann probability distribution $P(\mathcal{E}) = e^{-\mathcal{E}/kT}/kT$. The average value of the energy \mathcal{E} for this distribution is $\bar{\mathcal{E}} = kT$, which is the classical law of equipartition of energy. To calculate this value of $\bar{\mathcal{E}}$, we integrate $\mathcal{E}P(\mathcal{E})$ from zero to infinity. This is just the quantity that is being averaged, \mathcal{E} , multiplied by the relative probability $P(\mathcal{E})$ that the value of \mathcal{E} will be found in a measurement of the energy. *Bottom:* A plot of $\mathcal{E}P(\mathcal{E})$. The area under this curve gives the value of $\bar{\mathcal{E}}$.

Boltzmannverteilung – Diskret



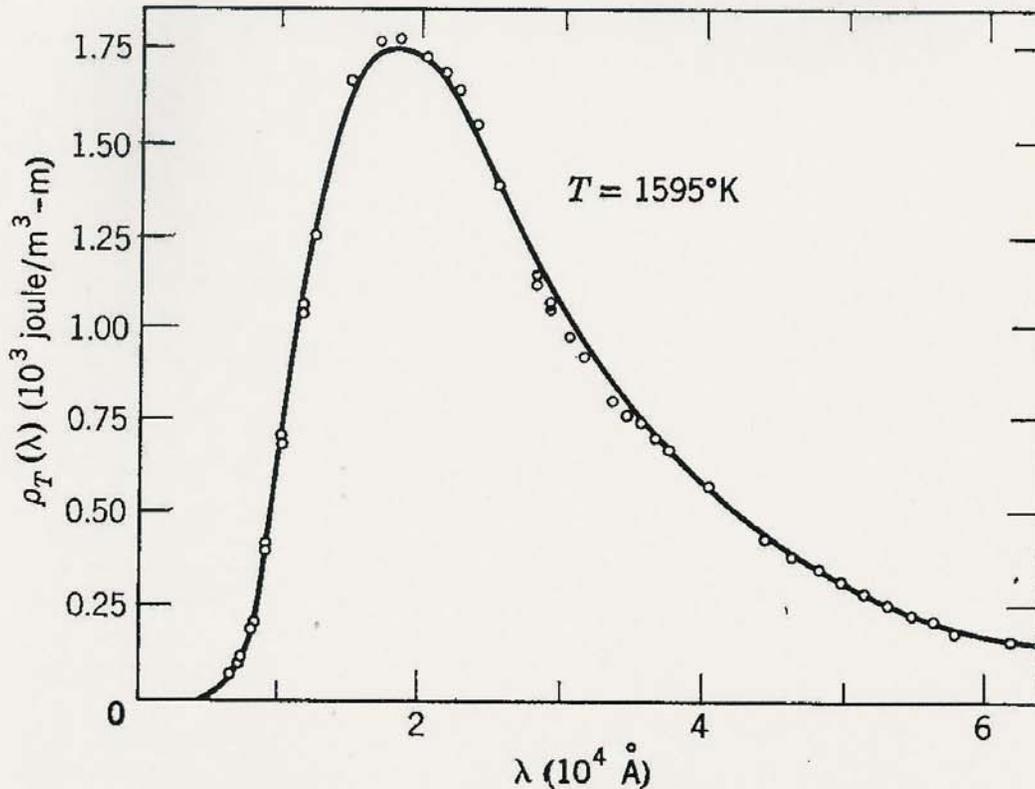
$$P(E_n) = \frac{1}{k_B T} e^{-E_n / k_B T} \quad n = 0, 1, 2, 3, \dots$$

$$E_n = n \cdot \Delta E$$

$$\bar{E} = \frac{\sum_{n=0}^{\infty} E_n \cdot P(E_n)}{\sum_{n=0}^{\infty} P(E_n)}$$

Figure 1-10 *Top:* If the energy ϵ is not a continuous variable but is instead restricted to discrete values $0, \Delta\epsilon, 2\Delta\epsilon, 3\Delta\epsilon, \dots$, as indicated by the ticks on the ϵ axis of the figure, the integral used to calculate the average value $\bar{\epsilon}$ must be replaced by a summation. The average value is thus a sum of areas of rectangles, each of width $\Delta\epsilon$, and with heights given by the allowed values of ϵ times $P(\epsilon)$ at the beginning of each interval. In this figure $\Delta\epsilon \ll kT$, and the allowed energies being closely spaced the area of all the rectangles differs but little from the area under the smooth curve. Thus the average value $\bar{\epsilon}$ is nearly equal to kT , the value found in Figure 1-9. *Middle:* $\Delta\epsilon \simeq kT$, and $\bar{\epsilon}$ has a smaller value than it has in the case of the top figure. *Bottom:* $\Delta\epsilon \gg kT$, and $\bar{\epsilon}$ is further reduced. In all three figures the rectangles show the contribution to the total area of $\epsilon P(\epsilon)$ for each allowed energy. The rectangle for $\epsilon = 0$ of course is always of zero height. This will make a large effect on the total area if the widths of the rectangles are large.

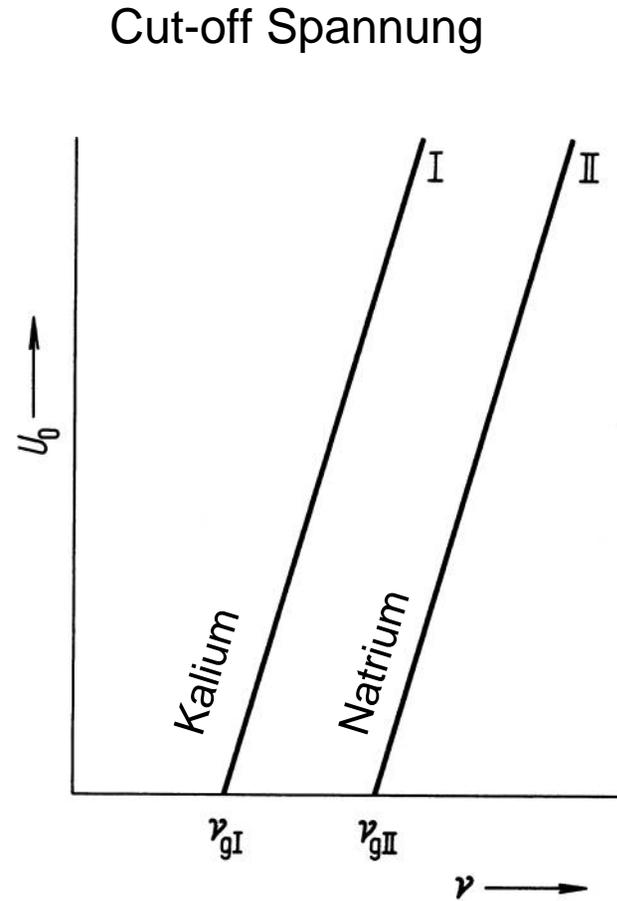
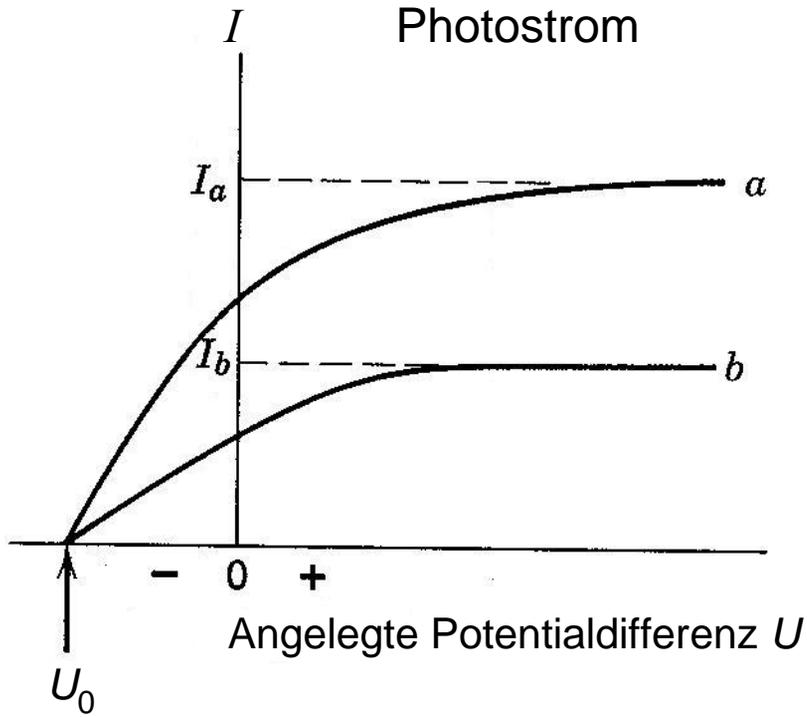
Plancksches Spektrum eines Schwarzen Körpers



$$\rho_T(\lambda) = \frac{8\pi hc}{\lambda^5} \frac{1}{e^{hc/\lambda k_B T} - 1}$$

Figure 1-11 Planck's energy density prediction (solid line) compared to the experimental results (circles) for the energy density of a blackbody. The data were reported by Coblentz in 1916 and apply to a temperature of 1595°K . The author remarked in his paper that after drawing the spectral energy curves resulting from his measurements, "owing to eye fatigue it was impossible for months thereafter to give attention to the reduction of the data." The data, when finally reduced, led to a value for Planck's constant of $6.57 \times 10^{-34} \text{ joule-sec}$.

Photoelektrischer Effekt – Experimente



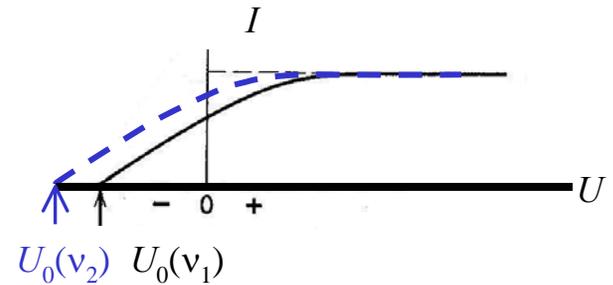
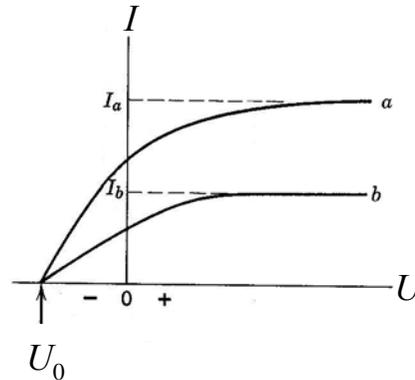
Photoelektr. Effekt – Widersprüche zw. Expt. und klassischer Wellen-Theorie des Lichts

1. $E_{\text{kin}} = eU_0$ sollte mit der Lichtintensität I_{licht} zunehmen:

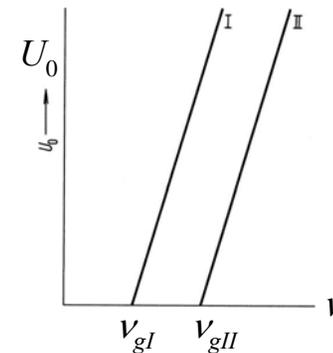
$$I_{\text{Licht}} \propto |\vec{E}|^2$$

$$\Rightarrow \vec{F}_E = e\vec{E} \propto I_{\text{Licht}}$$

$$\Rightarrow E_{\text{kin}} \propto I_{\text{Licht}}$$

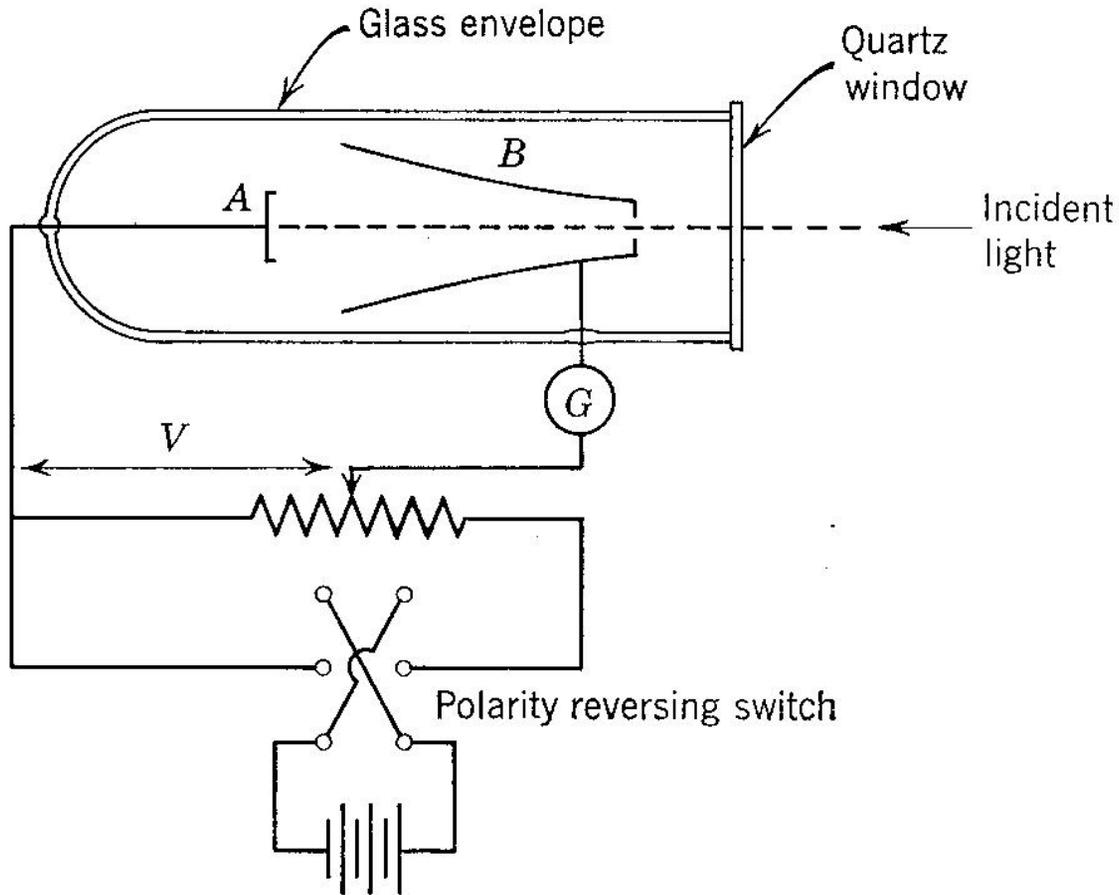


2. Photoelektr. Effekt sollte bei allen Frequenzen auftreten, sofern die Lichtintensität nur hinreichend hoch ist \Rightarrow keine „Cutoff“-Frequenz

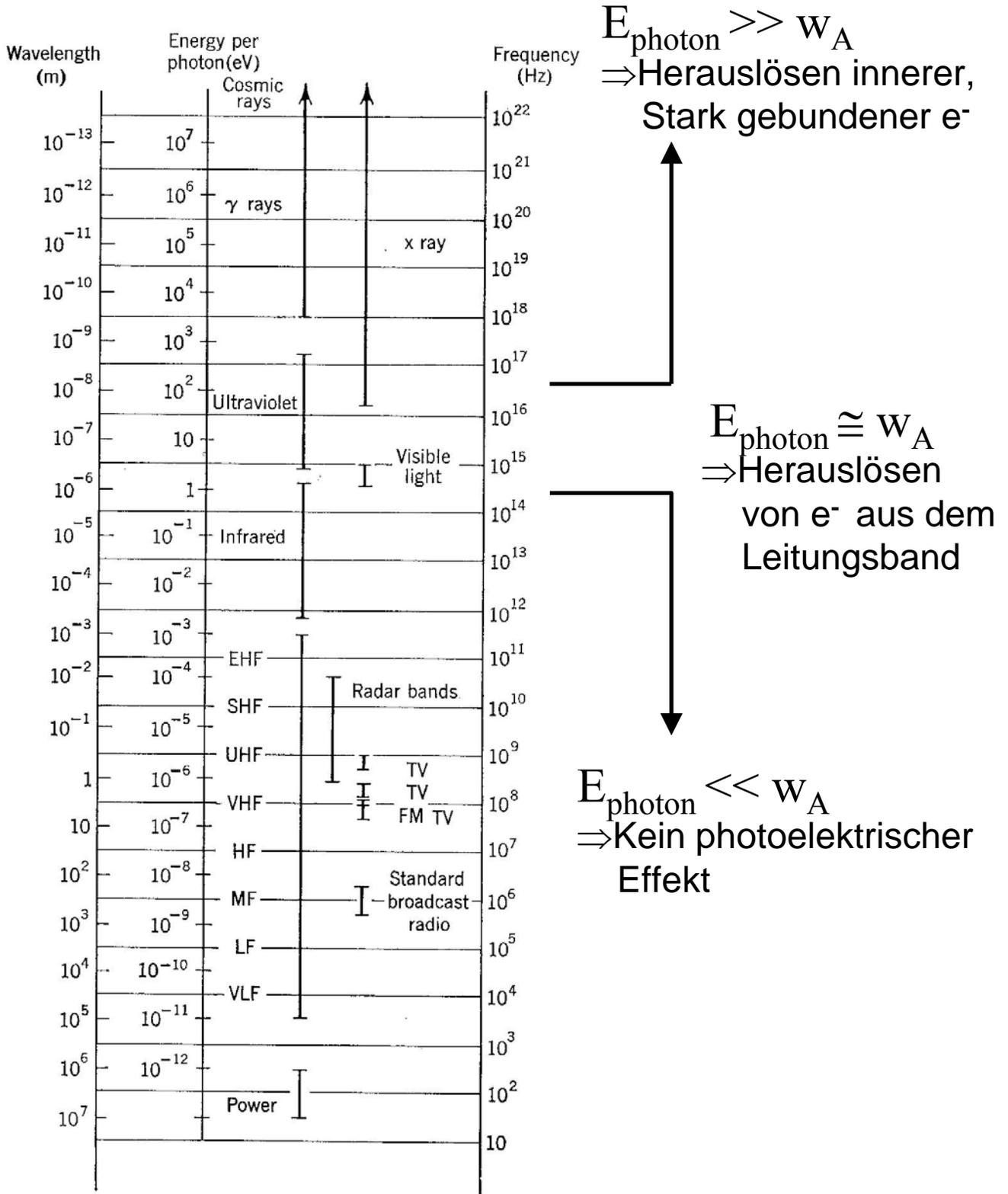


3. Zeitverzögerung zwischen Absorption und Photoelektronenemission – im Expt. nie beobachtet

Photoelektrischer Effekt – Messapparatur



Photoelektrischer Effekt und elektromagnetisches Spektrum



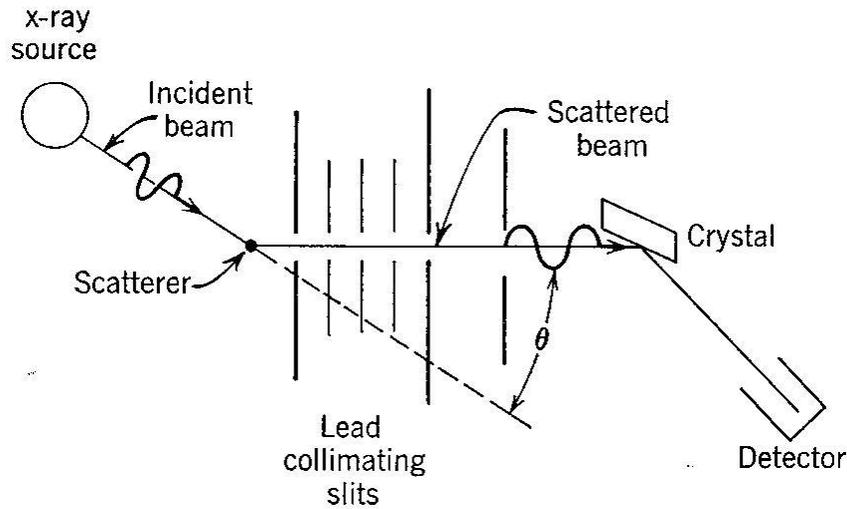
Photoelektrischer Effekt – Historie

$$h = 6.6262 \times 10^{-34} \text{ joule-sec}$$

To quote Millikan: “The photoelectric effect . . . furnishes a proof which is quite independent of the facts of blackbody radiation of the correctness of the fundamental assumption of the quantum theory, namely, the assumption of a discontinuous or explosive emission of the energy absorbed by the electronic constituents of atoms from . . . waves. It materializes, so to speak, the quantity h discovered by Planck through the study of blackbody radiation and gives us a confidence inspired by no other type of phenomenon that the primary physical conception underlying Planck’s work corresponds to reality.”

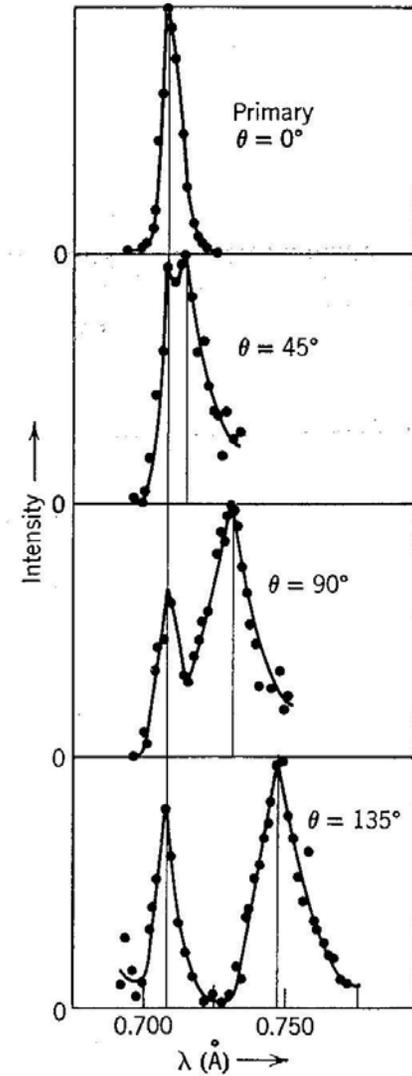
In 1921 Einstein received the Nobel Prize for predicting theoretically the law of the photoelectric effect. Before Millikan’s complete experimental validation of this law in 1914, Einstein was recommended to membership in the Prussian Academy of Sciences by Planck and others. Their early negative attitude toward the photon hypothesis is revealed in their signed affidavit, praising Einstein, in which they wrote: “Summing up, we may say that there is hardly one among the great problems, in which modern physics is so rich, to which Einstein has not made an important contribution. That he may have sometimes missed the target in his speculations, as, for example, in his hypothesis of light quanta (photons), cannot really be held too much against him, for it is not possible to introduce fundamentally new ideas, even in the most exact sciences, without occasionally taking a risk.”

Compton-Effekt



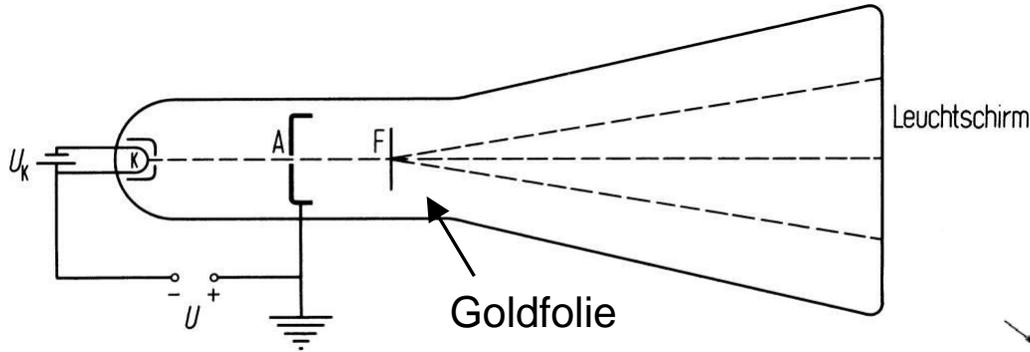
$$\lambda' = \lambda + \lambda_C (1 - \cos \theta)$$

$$\lambda_C = \frac{h}{m_e c} = 1.2 \cdot 10^{-12} \text{ m} \quad \text{Compton-Wellenlänge}$$

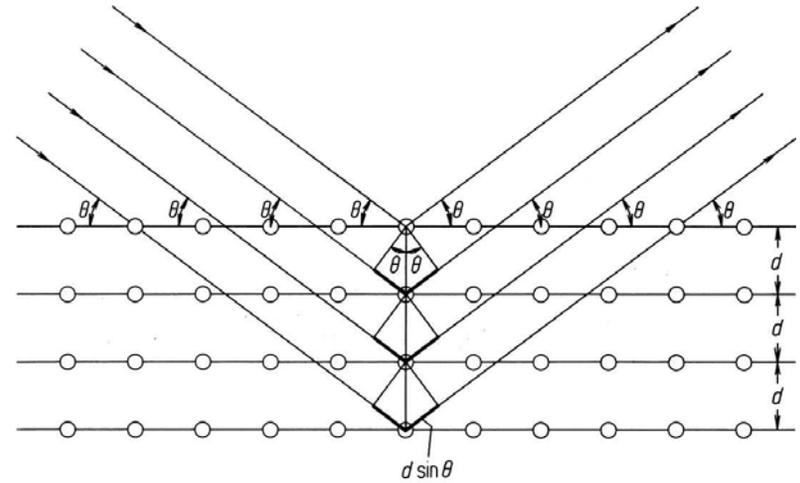


Wellennatur des Elektrons - Elektronenbeugung

G.P. Thomson 1927: Messapparatur



Bragg-Streuung



Braggsche Gleichung

$$n \cdot \lambda = 2d \sin \theta_n$$

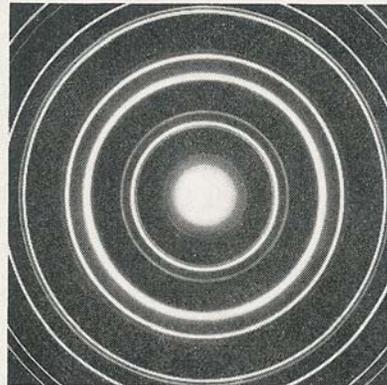
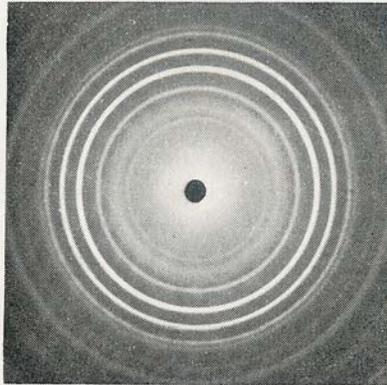
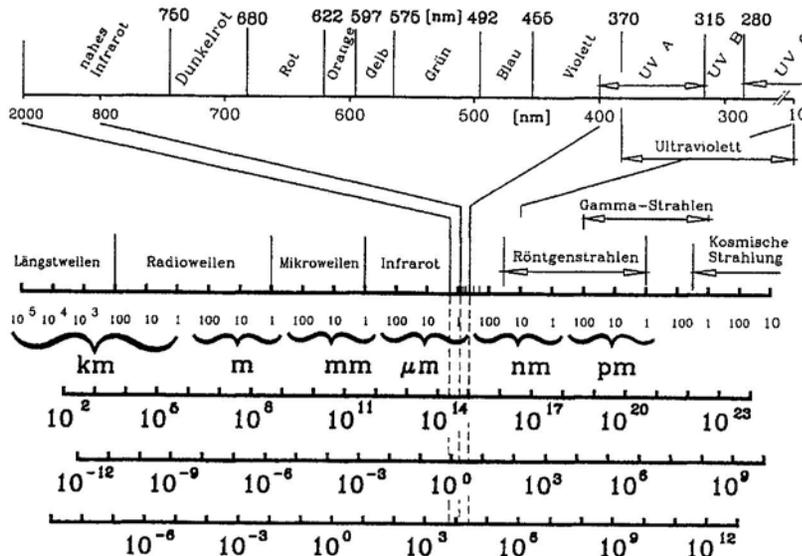
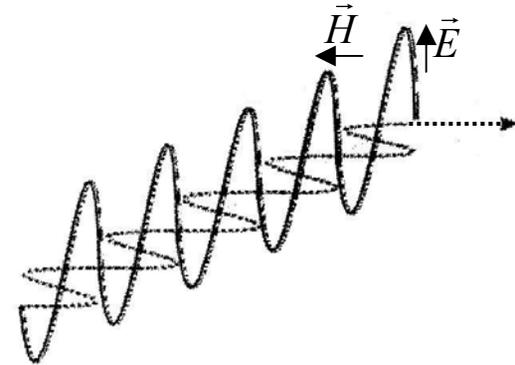
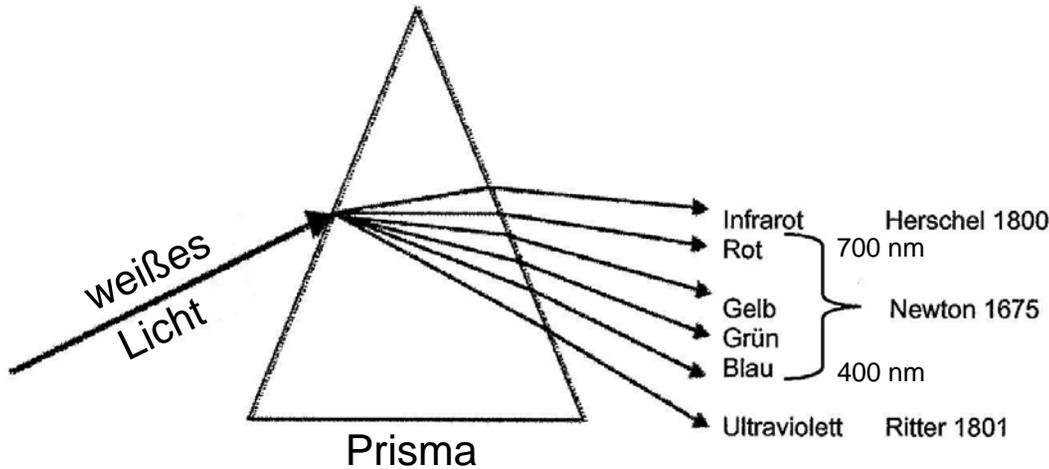


Figure 3-4 Top: The experimental arrangement for Debye-Scherrer diffraction of x rays or electrons by a polycrystalline material. Bottom left: Debye-Scherrer pattern of x-ray diffraction by zirconium oxide crystals. Bottom right: Debye-Scherrer pattern of electron diffraction by gold crystals.

Grundlagen der Spektroskopie



Sichtbares Licht

Elektromagnetisches
Spektrum des Lichts

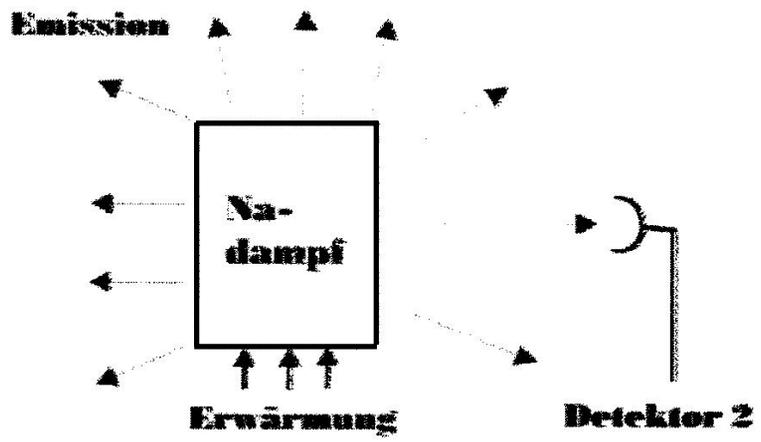
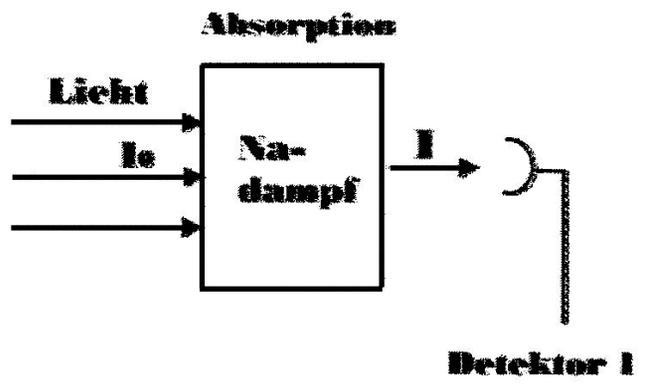
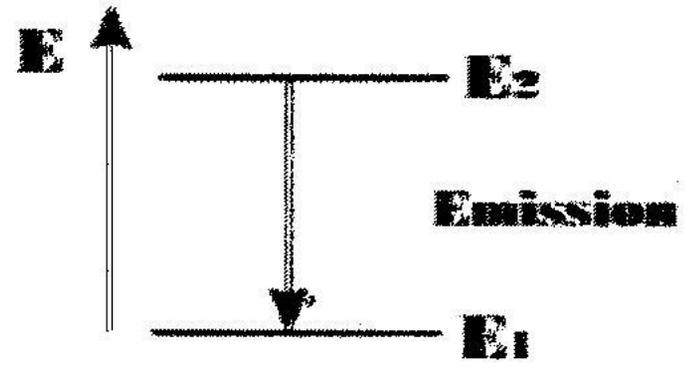
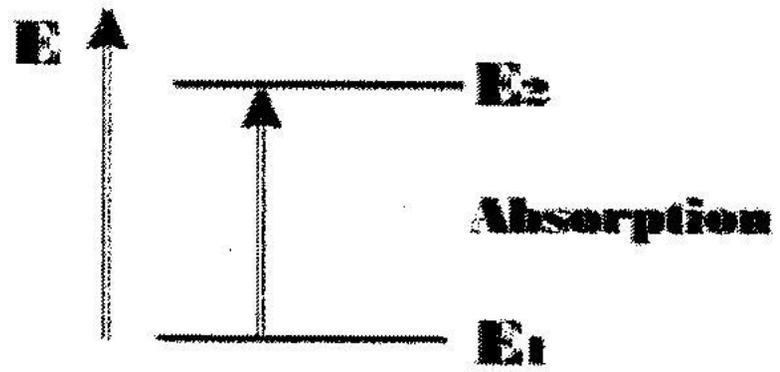
Wellenlänge, λ

Frequenz, ν in s^{-1}

Energie, E in eV

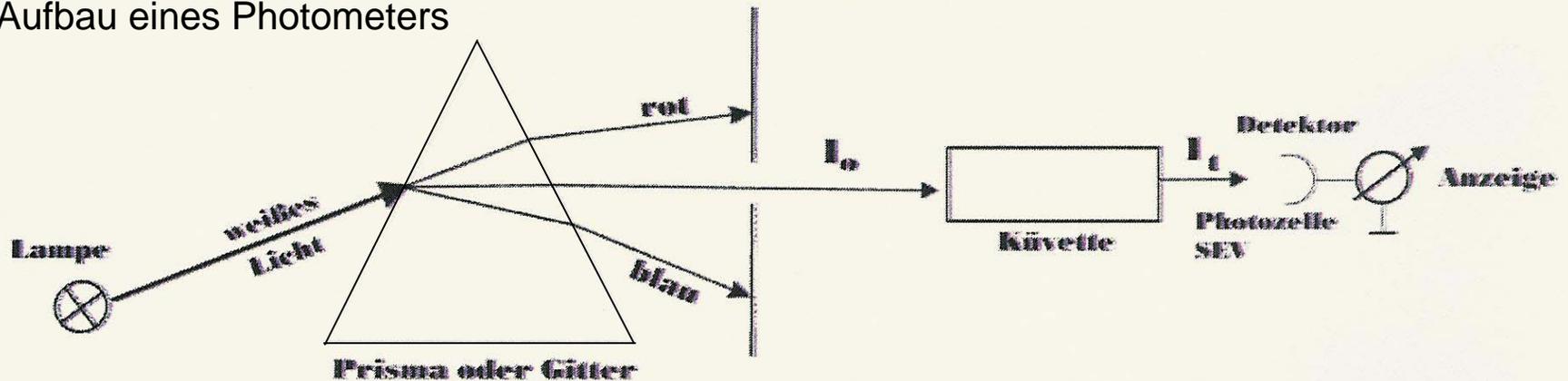
Wellenzahl, $\frac{1}{\lambda}$ in cm^{-1}

Absorption und Emission



Grundlagen der Spektroskopie

Aufbau eines Photometers



Mechanismen der Deaktivierung

